When does relativistic locality imply subsystem locality?

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introduction

two notions of locality

Relativistic



Spacetime regions



Quantum Information



Systems

introduction

low energy quantum gravity

If we detect gravity mediated entanglement, then gravity cannot be both:

classical

local

Spin Entanglement Witness for Quantum Gravity

Sougato Bose, Anupam Mazumdar, Gavin W. Morley, Hendrik Ulbricht, Marko Toroš, Mauro Paternostro, Andrew A. Geraci, Peter F. Barker, M. S. Kim, and Gerard Milburn

Phys. Rev. Lett. 119, 240401 – Published 13 December 2017



theory

Gravitationally Induced Entanglement between Two Massive Particles is Sufficient Evidence of Quantum Effects in Gravity constructor

C. Marletto and V. Vedral Phys. Rev. Lett. 119, 240402 - Published 13 December 2017

A no-go theorem on the nature of the gravitational field beyond quantum theory

Thomas D. Galley¹, Flaminia Giacomini¹, and John H. Selby²

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Published: Eprint: Doi: Citation:

2022-08-17, volume 6, page 779 arXiv:2012.01441v7 https://doi.org/10.22331/q-2022-08-17-779 Quantum 6, 779 (2022).

GPTs





introduction

a foundation for subsystem locality?

- Subsystem locality is a often good operational assumption.
 - Axiomatically assumed in quantum foundations works
 - Motivated by relativistic intuition
- **Does subsystem locality hold in QFT?**
 - **Does a field mediate interactions, in the QI sense?** \bullet
 - To what extent does this property hold in nature?





plan

- two false starts
- the scalar field theory case
- open questions

two false starts

two false starts

Suzuki-Trotter

$H = H_A + H_B + H_C + H_{AC} + H_{BC}$

two false starts **Suzuki-Trotter** $H = H_{AC} + H_{RC}$ $\implies U(t) = e^{-i(H_{AC} + H_{BC})t} \neq e^{-iH_{AC}t}e^{-iH_{BC}t}$ $= \lim \left(e^{-iH_{AC}t/n} e^{-iH_{BC}t/n} \right)^n$ $n \rightarrow \infty$

Arbitrarily good approximation

but no input from relativity.

two false starts **QED in Coulomb gauge** $H = H_1 + H_2 + H_{A_{\perp}}^{\text{rad}} + \frac{q_1 q_2}{|\mathbf{x}_1 - \mathbf{x}_2|} - \int d^3 \mathbf{x} A_{\perp}(\mathbf{x}) \cdot (J_1(\mathbf{x}) + J_2(\mathbf{x}))$

If two particles are at rest, and there is no radiation, then

$$H \approx \frac{4}{|\mathbf{x}_1|}$$

Not subsystem local!

 $\frac{q_1 q_2}{|-\mathbf{x}_2|}$

massive scalar field

massive scalar field a positive result

Concrete example:

- Two particles coupled to a massive scalar field, in a specific regime.
- **Evolution is subsystem local, up to some phases.**
- Microcausality ($[\hat{\phi}(x), \hat{\phi}(x')] = 0$ if x, x' spacelike) eliminates the phases.
 - **Relativistic locality implies subsystem locality.**

arXiv:2305.05645 (quant-ph)

[Submitted on 9 May 2023] Relativistic locality can imply subsystem locality

Andrea Di Biagio, Richard Howl, Časlav Brukner, Carlo Rovelli, Marios Christodoulou

massive scalar field three key assumptions

- Support of the matter wavefunctions contained within two distinct spacetime regions.
- Matter in quantum-controlled superposition of semi-classical states.





derivation sketch

derivation setup

quantum-controlled dynamics

$$\hat{H}_{A}(t) = \sum_{r} |r X r| \otimes \hat{H}_{A}^{r}(t)$$

$$\hat{H}_{\text{int}} = \int d^3 \mathbf{x} \, \hat{\phi}(\mathbf{x})$$

 $\hat{H}(t) = \hat{H}_A(t) + \hat{H}_B(t) + \hat{H}_0 + \hat{H}_{int}$

kinetic field term

$$\hat{H}_0 = \int \frac{\mathrm{d}^3 \mathbf{k}}{(2\pi)^3} \,\omega_{\mathbf{k}} \hat{a}_{\mathbf{k}}^{\dagger} \hat{a}_{\mathbf{k}}$$

local interaction

 $(\hat{\mu}_A(\mathbf{x}) + \hat{\mu}_B(\mathbf{x}))$

derivation qudit-controlled dynamics

No back action on the qudits + matter in superposition of pointer states:

$$|\Psi(t)\rangle = \sum_{rs} c_{rs} |rs\rangle|_{rs}$$

articles:
$$\frac{\mathrm{d}}{\mathrm{d}t} |\psi_A^r(t)\rangle = -i\hat{H}_A^r(t) |\psi_A^r(t)|$$

ield:
$$\frac{\mathrm{d}}{\mathrm{d}t} |\phi^{rs}(t)\rangle = -i(\hat{H}_0 + \hat{H}_2)$$

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 $|\psi_A^r(t)\rangle |\psi_B^s(t)\rangle |\phi^{rs}(t)\rangle$

$(t)\rangle$

$\hat{H}_{int}^{rs}(t) \left| \phi^{rs}(t) \right\rangle \qquad \hat{H}_{int}^{rs}(t) = \left\langle \psi_A^r(t) \psi_B^s(t) \left| \hat{H}_{int} \right| \psi_A^r(t) \psi_B^s(t) \right\rangle$

Evolution of the whole system: $\hat{U} = \sum |rs | rs | \otimes \hat{U}_A^r \otimes \hat{U}_B^s \otimes \hat{U}_{\phi}^{rs}$ rs



derivation condition for subsystem locality

$\hat{U} = \sum_{A} |rs| \langle rs| \otimes \hat{U}_{A}^{r} \otimes \hat{U}_{B}^{s} \otimes \hat{U}_{d}^{rs}$ rs

but if we had $\forall rs$: $\hat{U}^{rs}_{\phi} = \hat{U}^{r}_{\phi} \circ \hat{U}^{s}_{\phi}$

 $\hat{U} = \left(\sum_{s} |s\rangle\langle s| \otimes \hat{U}_{B}^{s} \otimes \hat{U}_{\phi}^{s}\right) \circ \left(\sum_{r} |r\rangle\langle r| \otimes \hat{U}_{A}^{r} \otimes \hat{U}_{\phi}^{r}\right)$





derivation evolution of the field



 $\hat{U} = \sum |rs| \langle rs| \otimes \hat{U}_A^r \otimes \hat{U}_B^s \otimes \hat{U}_\phi^{rs}$ rs

 $\frac{\mathrm{d}}{\mathrm{d}t}\hat{U}_{\phi}^{rs}(t) = -i\hat{H}^{rs}(t)\hat{U}_{\phi}^{rs}(t)$



derivation evolution of the field



quantum field with classical source!

 $\frac{\mathrm{d}}{\mathrm{d}t}\hat{U}_{\phi}^{rs}(t) = -i\hat{H}^{rs}(t)\hat{U}_{\phi}^{rs}(t)$

 $\hat{H}^{rs}(t) = \hat{H}_0 + \langle \psi_A^r(t)\psi_B^s(t) | \hat{H}_{int} | \psi_A^r(t)\psi_B^s(t) \rangle$

exact solution

 $\hat{U}_{\phi}^{rs} = e^{i\Omega^{rs}} \hat{D}^{rs} e^{-i\hat{H}_0(t_2 - t_1)}$



derivation subsystem locality?

$$\hat{U}_{\phi}^{rs} = e^{i\Omega^{rs}}\hat{D}^{rs}e^{-i\hat{H}_0(t_2-t_1)}$$

$$\hat{U} = \sum_{sr} e^{i\tilde{\Omega}^{rs}} \left(|s\rangle\langle s| \hat{U}_B^s \otimes \hat{U}_\phi^s \right) \circ \left(|r\rangle\langle r| \hat{U}_A^r \otimes \hat{U}_\phi^r \right) \circ e^{-i\hat{H}_0(t_2 - t_1)}$$

almost subsystem local!

 $e^{i\tilde{\Omega}^{rs}}\hat{U}^{r}_{\phi}\hat{U}^{s}_{\phi}e^{-i\hat{H}_{0}(t_{2}-t_{1})}$

full evolution:

—

derivation the phase

$$\tilde{\boldsymbol{\Omega}}^{rs} = -i \iint_{t_1}^{t_2} \mathrm{d}t \mathrm{d}t' \iint \mathrm{d}^3 \mathbf{x} \mathrm{d}^3 \mathbf{x}' \,\mu_A^r(t, \mathbf{x}) \mu$$
$$-i \int_{t_1}^{t_2} \mathrm{d}t \int_{t_1}^t \mathrm{d}t' \iint \mathrm{d}^3 \mathbf{x} \mathrm{d}^3 \mathbf{x}' \big(\mu_A^r(t, \mathbf{x}) \big) \mathrm{d}^3 \mathbf{x} \mathrm{d}^3 \mathbf{x}' \big(\mathbf{x} \mathrm{d}^3 \mathbf{x}' \big(\mathbf{x} \mathrm{d}^3 \mathbf{x}' \big) \big(\mathbf{x} \mathrm{d}^3 \mathbf{x} \mathrm{d}^3 \mathbf{x}' \big) \mathrm{d}^3 \mathbf{x} \mathrm{d}^3 \mathbf{x}' \big(\mathbf{x} \mathrm{d}^3 \mathbf{x} \mathrm{d}^3 \mathbf{x}' \big(\mathbf{x} \mathrm{d}^3 \mathbf{x}' \big) \mathrm{d}^3 \mathbf{x} \mathrm{d}^3 \mathbf{x}' \big(\mathbf{x} \mathrm{d}^3 \mathbf{x} \mathrm{d}^3 \mathbf{x}' \big(\mathbf{x} \mathrm{d}^3 \mathbf{x} \mathrm{d}^3 \mathbf{x}' \big) \mathrm{d}^3 \mathbf{x} \mathrm{d}^3 \mathbf{x}' \big(\mathbf{x} \mathrm{d}^3 \mathbf{x} \mathrm{d}^3 \mathbf{x}' \big(\mathbf{x} \mathrm{d}^3 \mathbf{x} \mathrm{d}^3 \mathbf{x}' \big(\mathbf{x} \mathrm{d}^3 \mathbf{x} \mathrm{d}^3 \mathbf{x}' \big) \mathrm{d}^3 \mathbf{x} \mathrm{d}^3 \mathbf{x}' \big(\mathbf{x} \mathrm{d}^3 \mathbf{x} \mathrm{d}^3 \mathbf{x}' \big) \mathrm{d}^3 \mathbf{x} \mathrm{d}^3 \mathbf{x}' \big(\mathbf{x} \mathrm{d}^3 \mathbf{x}' \big(\mathbf{x} \mathrm{d}^3 \mathbf{x} \mathrm{d}^3 \mathbf{x}' \big(\mathbf{x} \mathrm{d}^3 \mathbf{x} \mathrm{d}^3 \mathbf{x}' \big(\mathbf{x} \mathrm{d}^3 \mathbf{x}' \big(\mathbf{x} \mathrm{d}^3 \mathbf{x} \mathrm{d}^3 \mathbf{$$

 $u_B^{S}(t',\mathbf{X}')[\hat{\phi}_I(t,\mathbf{X}),\hat{\phi}_I(t',\mathbf{X}')]$

$(1)\mu_B^s(t',\mathbf{x}') + \mu_B^r(t,\mathbf{x})\mu_A^s(t',\mathbf{x}'))[\hat{\phi}_I(t,\mathbf{x}),\hat{\phi}_I(t',\mathbf{x}')]$



derivation relativistic locality

 $\tilde{\Omega}^{rs} = -i \iint_{\tau}^{\iota_2} \mathrm{d}^4 x \mathrm{d}^4 x \,\mu_A^r(x) \mu_B^s(x') [\hat{\phi}_I(x), \hat{\phi}_I(x')] - \cdots$

Microcausality:

$[\hat{\phi}_I(x), \hat{\phi}_I(x')] = 0$

if x and x' are spacelike

derivation relativistic locality

$$\tilde{\Omega}^{rs} = -i \iint_{t_1}^{t_2} \mathrm{d}^4 x \mathrm{d}^4 x \,\mu_A^r(x) \mu_B^s(x') [\hat{\phi}_I(x)]$$

if $\operatorname{supp} \mu_A^r$, $\operatorname{supp} \mu_B^s$ are spacelike

then $\tilde{\Omega}^{rs} = 0$





X

derivation relativistic locality

$$\tilde{\Omega}^{rs} = -i \iint_{t_1}^{t_2} \mathrm{d}^4 x \mathrm{d}^4 x \,\mu_A^r(x) \mu_B^s(x') [\hat{\phi}_I(x')] = -i \int_{t_1}^{t_2} \mathrm{d}^4 x \,\mathrm{d}^4 x \,\mathrm{d}^$$

if $\operatorname{supp} \mu_A^r$, $\operatorname{supp} \mu_B^s$ are spacelike $\forall rs$

then $\tilde{\Omega}^{rs} = 0$ $\forall rs$







relativistic locality implies subsystem locality

 $\hat{U}(t_1, t_2) = \hat{U}_{B\phi} \circ \hat{U}_{A\phi} \circ e^{-i\hat{H}_0(t_2 - t_1)}$ then





conclusion

Summary

- simple model, and only in a certain approximate regime.
- This is to be expected.
- Interesting intersection for RQI.

Subsystem locality and relativistic locality are related but different notions.

Relativistic locality (via microcausality) implies subsystem locality, in a



conclusion

Open questions

- To what extend can this result be generalised?
 - Massless, gauge fields?
- Perhaps we need to leverage tools from AQFT.
- If result cannot be generalised, what is the impact on quantum foundations?

thank you